

Nuclear Reactor Physics Lecture Note (11)
-Neutron Spectrum(2) Resonance Absorption-

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9.2 Resonance Absorption

(1) Resonance Escape Probability

Resonance escape probability p

Fraction of fission neutrons that manage to slow down from fission to thermal energies without being absorbed

Resonance escape probability in hydrogen

Assumptions

Infinite medium of hydrogen

An infinitely massive absorber distributed uniformly

All neutron slowing down due to elastic scattering with hydrogen (ignore inelastic scattering)

Neutron source (neutron energy E_s) uniformly distributed in the medium

The resonance escape probability to energy E , $p(E)$, is given by

$$p(E) = \exp \left[- \int_E^{E_s} \frac{dE' \Sigma_a(E')}{E' \Sigma_t(E')} \right] \quad \dots (1)$$

where,

$\Sigma_a(E')$: macroscopic absorption cross section at energy E'

$\Sigma_t(E')$: macroscopic total cross section at energy E'

(2) Breit-Wigner Single-Level Resonance Formula and Doppler Effect

(a) Breit-Wigner single-level resonance formula

(Energy dependence of absorption cross section for resonance which are spaced widely apart)

$$\sigma_\gamma(E_c) = \sigma_0 \frac{\Gamma_\gamma}{\Gamma} \left(\frac{E_0}{E_c} \right)^{\frac{1}{2}} \left[4 \left(\frac{E_c - E_0}{\Gamma} \right)^2 + 1 \right]^{-1} \quad \dots (2)$$

where,

E_0 : the energy at which resonance occurs

Γ : total time width of the resonance

(characterizing the width of the energy level and the full width at half-maximum (FWHM) of the resonance)

Γ_γ : radiation line width

(characterizing the probability that the compound nucleus will decay via gamma emission)

σ_0 : total cross section at the resonance energy E_0

E_c : energy in center of mass system ($E_c = \frac{M}{m + M} E$)

- For low energy ($E \ll E_0$) $\sigma_\gamma \propto \frac{1}{E^2}$ (or $\frac{1}{v}$)
- For high energy ($E_0 \ll E$) $\sigma_\gamma \propto \frac{1}{E^2}$

(b) Doppler effect on cross section resonance

The nuclei are in thermal motion.

The relative speed between the nucleus and the neutron can be either greater or less than neutron speed.

→ Doppler shift effect

as temperature increase, the resonance broadens, while its peak magnitude decreases.

Doppler-broadened Breit-Wigner resonance cross section (Bethe-Placzek approximation)

When $E \sim E_0$

$$\sigma_\gamma(E) = \sigma_0 \frac{\Gamma_\gamma}{\Gamma} \psi(\zeta, \chi) \quad \dots (3)$$

$$\zeta = \Gamma \left(\frac{A}{4E_0 kT} \right)^{\frac{1}{2}}$$

$$\chi = 2 \left(\frac{E - E_0}{\Gamma} \right) \quad T : \text{absolute temperature, } k : \text{boltsmann constant}$$

(3) Resonance escape probability in infinite dilution approximation p^∞

Assumptions

Infinite hydrogen medium

Ignoring absorption by hydrogen and scattering by absorbing nucleus

Absorbing nucleus are infinitely diluted

$$p^\infty = \exp \left[- \frac{\pi N_A \sigma_0 \Gamma_\gamma}{2 N_H \sigma_S^H E_0} \right] \quad \dots (4)$$

where,

N_A : number density of absorbing nucleus

N_H : number density of hydrogen

σ_s^H : microscopic scattering cross section of hydrogen (assumed as constant)

Feature of p^∞

No temperature dependence

Increases with increasing moderator (hydrogen) density

Ex : Moderator void in light water reactor

→low moderator density

→low p

→negative reactivity

Decreases in low E_0

→lower energy resonances give a lot of effect on reactivity.

Ex : 6.67eV resonance in ^{238}U