

image theory. when image theory is applied, region I is substituted by free space and the equivalent magnetic current becomes  $2\mathbf{M}_e$ . this argument is valid only for  $z > 0$ , where the equivalent electromagnetic field becomes  $2\mathbf{E}_{em}$  and  $2\mathbf{H}_{em}$  similar to the equations (7)(8).

$$\mathbf{E}_e = 2\mathbf{E}_{em} = \hat{\mathbf{x}}E_0 \exp(-jk_0z) = \mathbf{E}_o \quad (15)$$

$$\mathbf{H}_e = 2\mathbf{H}_{em} = \hat{\mathbf{y}} \frac{E_0}{Z_0} \exp(-jk_0z) = \mathbf{H}_o \quad (16)$$

Therefore, the original field appears in region II and is observed in fig.7. In region I the electromagnetic field vanishes because of PEC.

#### 1.4 Equivalent model for region I when region II is perfect electric conductor (PEC)

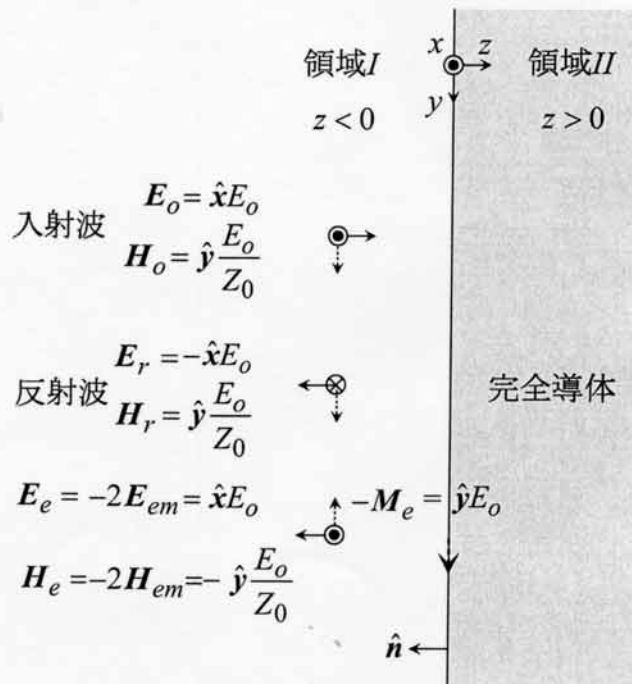


Figure 9: 領域 II を完全導体で満たした領域 I の等価モデル (1)

Figure 9 shows case when region II is PEC. In this case when the original field ( $\mathbf{E}_o, \mathbf{H}_o$ ) hits the electric wall, the reflected field ( $\mathbf{E}_r, \mathbf{H}_r$ ) is induced. On the electric wall the equivalent magnetic current  $-\mathbf{M}_e$  produce the equivalent electromagnetic field ( $\mathbf{E}_e, \mathbf{H}_e$ ). For the total contribution the three components are needed. The reflected wave ( $\mathbf{E}_r, \mathbf{H}_r$ ) is defined so that the tangential components of electric field vanishes on the electric wall  $S$ .

$$\mathbf{E}_r = -\hat{\mathbf{x}}E_0 \exp(jk_0z) \quad (17)$$

$$\mathbf{H}_r = \hat{\mathbf{y}} \frac{E_0}{Z_0} \exp(jk_0z) \quad (18)$$

Because the electric wall  $S$  is infinite plane, the electromagnetic field ( $\mathbf{E}_e, \mathbf{H}_e$ ) produced by equivalent magnetic current  $-\mathbf{M}_e$  for  $z < 0$  might be calculated using image theory. When